# Solving (3+1)-dimensional Jimpo-Miwa equations using Riccati mapping equation method 

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#### Abstract

Several new exact solutions of the $(3+1)$-dimensional Jimbo-Miwa equation and its two extended forms have been successfully obtained by using the Riccati mapping equation method. Although some of the dispersion relations are the same as in [17] our solutions consedered new .


Key words: Riccati equation mapping method, (3+1)-dimensional Jimpo-Miwa equations, Exact solutions.
الملخص الكثير من الحلول التامة الجديدة لمعادلة جامبو -ميوا ذات البعد (3 + 1 ) وصيغتاها الموسعتان تم الحصول عليها بنجاح باستخدام طريقة معادلة ريكاتي. بالرغم من ان علاقات التثتت تعتبر نفسها في [17] الا ان حلولنا تعتبر
الكلمات المفتاحية: طريقة معادلة ريكاتي، معالادت جمبو -ميوا ذات البعد (3+1)، حلول تامة

## 1 Introduction

Various powerful methods have been used to refine various types of solutions for nonlinear partial differential equations models that are used to explain a broad variety of physical phenomena. . Some of these methods are $\tan \left(F\left(\frac{5}{2}\right)\right)$-expansion method [1], the sine-cosine function method [2], the variational iteration method [3,4], Riccati mapping equation method [5-7], the exp-function method [8,9], the generalized and the improved generalized tanhcoth method [10,11], the new modified simple equation method [12], the simplified Hirota's direct method $[13,14]$, tanh method, sine-Gordon expansion method, and the extended sinhGordon equation expansion method $[15,16]$. Wazwaz [17] has obtained multiple soliton solutions of distinct physical structures for $(3+1)$-dimensional Jimbo-Miwa and its two extended equations by using simplified Hirota's method. Many other solutions have been found by several researchers [18-20].
In this article, we use Riccati mapping equation method to extract many new solutions of both the $(3+1)$-dimensional Jimbo-Miwa equation and its two extended forms as well. The dispersion relations are found to be distinct for each equation.

## 2 Description of the Riccati equation mapping method

Consider a given nonlinear partial differential equation with the independent variable $X=\left(t, x_{1}, x_{2}, \ldots, x_{n}\right)$ and the dependent variable $u(X)$ in the following form;

$$
\begin{equation*}
W\left(u, u_{t}, u_{x_{i}}, u_{x_{i} x_{i}}, u_{x_{i} x_{j}}, \ldots\right)=0 \tag{1}
\end{equation*}
$$

where $W$ is a general polynomial function of its arguments. In order to solve Eq.(1) by using the proposed method, the following main steps are given;

Step 1: Substituting

$$
\begin{equation*}
u=u(\zeta), \quad \zeta=k x+r y+s z-c t \tag{2}
\end{equation*}
$$

into (1) yields an ordinary differential equation in $\zeta$ of the form

$$
\begin{equation*}
E\left(u, u^{\prime}(\zeta), u^{\prime \prime}(\zeta), u^{\prime \prime \prime}(\zeta), \ldots\right)=0 \tag{3}
\end{equation*}
$$

where $E$ is a general polynomial.
Step 2: Assuming that (3) has the formal solution

$$
\begin{equation*}
u(\zeta)=a_{0}+\sum_{i=1}^{m}\left(a_{i} \phi^{i}(\zeta)+b_{i} \phi^{i-1}(\zeta) \sqrt{\left(\sigma+\phi^{2}(\zeta)\right)}\right) \tag{4}
\end{equation*}
$$

with $\phi$ satisfies the Riccati equation;

$$
\begin{equation*}
\phi^{\prime}(\zeta)=\sigma+\phi^{2}(\zeta) \tag{5}
\end{equation*}
$$

where $\sigma, a_{0}, a_{i}$, and $b_{i}$ are constants to be determined later, while $m$ is a positive integer, which is called the balance number.

Step 3: Determining the positive integer $m$ by balancing the highest order derivatives and the nonlinear terms in Eq. (3).

Step 4: Substituting (4) along with Eq. (5) into (3) and collecting all the coefficients of
$\phi^{n}\left(\sqrt{\sigma+\phi^{2}(\zeta)}\right)^{m}, n=0, \pm 1, \pm 2, \pm 3, \ldots . . ; m=0, \pm 1$ then setting each coefficient to zero,
a set of algebraic equations is obtained for $a_{0}, a_{i}, b_{i}$, and $c$.
Step 5: Solving the system of algebraic equations in Step4 using Matlab or Mathematica software to find the values of $a_{0}, a_{i}, b_{i}$, and $c$.

Step 6: As Eq.(5) possesses the solutions;

$$
\phi(\zeta)=\left(\begin{array}{ll}
-\sqrt{-\sigma} \tanh (\sqrt{-\sigma} \zeta) & , \sigma<0  \tag{6}\\
-\sqrt{-\sigma} \operatorname{coth}(\sqrt{-\sigma} \zeta) & , \sigma<0 \\
\sqrt{\sigma} \tan (\sqrt{\sigma} \zeta) & , \sigma>0 \\
\sqrt{\sigma} \cot (\sqrt{\sigma} \zeta) & , \sigma>0 \\
-\frac{1}{\zeta} & , \sigma=0
\end{array}\right.
$$

substituting $a_{0}, a_{i}, b_{i}$, and (6) into (4) to obtain the exact solutions of Eq. (1).

## 3. Applications

In this section, new exact solutions of the nonlinear $(3+1)$-dimensional Jimpo-Miwa equation and its two extracted forms have been developed by using the described Riccati equation mapping method. Without loss of generality we will assume through out this paper that the integration constants are set to be zeros.

### 3.1 The exact solutions of $(3+1)$-dimensional Jimpo-Miwa equation:

$$
\begin{equation*}
u_{x x x y}+3 u_{y} u_{x x}+3 u_{x} u_{x y}+2 u_{y t}-3 u_{x z}=0 \tag{7}
\end{equation*}
$$

Using the wave transformations $u=u(\zeta), \quad \zeta=k x+r y+s z-c t$, where $k, r, s$ and $c$ are constants in equation (7) and integrating we get

$$
\begin{equation*}
k^{3} r u^{\prime \prime \prime}+3 k^{2} r\left(u^{\prime}\right)^{2}-(2 r c+3 k s) u^{\prime}=0 . \tag{8}
\end{equation*}
$$

Balancing $u^{\prime \prime \prime}$ with $\left(u^{\prime}\right)^{2}$ we get $m=1$. From (4) we have

$$
\begin{equation*}
u(\zeta)=a_{0}+a_{1} \phi(\zeta)+b_{1} \sqrt{\sigma+\phi^{2}(\zeta)} \tag{9}
\end{equation*}
$$

Substituting (9) and (5) into equation (8) and collecting all the coefficients of $\phi^{n}\left(\sqrt{\sigma+\phi^{2}(\zeta)}\right)^{m}$ and setting them equal to zero we obtain

Case 1: $a_{0}=0, a_{1}=-k, \quad b_{1}=-k$, and $c=-\left(\frac{k^{3} r \sigma+3 k s}{2 r}\right)$. The exact solution in this case is

$$
u(\zeta)=\left(\begin{array}{ll}
k \sqrt{-\sigma} \tanh (\sqrt{-\sigma} \zeta)-k \sqrt{\sigma} \operatorname{sech}(\sqrt{-\sigma} \zeta) & , \sigma<0  \tag{10}\\
k \sqrt{-\sigma} \operatorname{coth}(\sqrt{-\sigma} \zeta)+k \sqrt{\sigma} \operatorname{csch}(\sqrt{-\sigma} \zeta) & , \sigma<0 \\
-k \sqrt{\sigma} \tan (\sqrt{\sigma} \zeta)-k \sqrt{\sigma} \sec (\sqrt{\sigma} \zeta) & , \sigma>0 \\
-k \sqrt{\sigma} \cot (\sqrt{\sigma} \zeta)-k \sqrt{\sigma} \csc (\sqrt{\sigma} \zeta) & , \sigma>0 \\
0 & , \sigma=0
\end{array}\right.
$$

Remark: For $\sigma=-1$ the dispersion relation is completely agree with [17]
Case 2: $a_{0}=0, \quad a_{1}=-2 k, \quad b_{1}=0, \quad c=-\left(\frac{4 k^{3} r \sigma+3 k s}{2 r}\right)$ and the exact solution is

$$
u(\zeta)=\left(\begin{array}{ll}
2 k \sqrt{-\sigma} \tanh (\sqrt{-\sigma} \zeta) & , \sigma<0  \tag{11}\\
2 k \sqrt{-\sigma} \operatorname{coth}(\sqrt{-\sigma} \zeta) & , \sigma<0 \\
-2 k \sqrt{\sigma} \tan (\sqrt{\sigma} \zeta) & , \sigma>0 \\
-2 k \sqrt{\sigma} \cot (\sqrt{\sigma} \zeta) & , \sigma>0 \\
\frac{2 k}{\zeta} & , \sigma=0
\end{array}\right.
$$

Remark: For $\sigma=-\frac{1}{4}$ the dispersion relation is completely agree with [17]
Case 3: $a_{0}=0, \quad a_{1}=-k, \quad b_{1}=k, \quad c=-\left(\frac{k^{3} r \sigma+3 k s}{2 r}\right)$ and the exact solution is

$$
u(\zeta)=\left(\begin{array}{ll}
k \sqrt{-\sigma} \tanh (\sqrt{-\sigma} \zeta)+k \sqrt{\sigma} \operatorname{sech}(\sqrt{-\sigma} \zeta) & , \sigma<0  \tag{12}\\
k \sqrt{-\sigma} \operatorname{coth}(\sqrt{-\sigma} \zeta)-k \sqrt{\sigma} \operatorname{csch}(\sqrt{-\sigma} \zeta) & , \sigma<0 \\
-k \sqrt{\sigma} \tan (\sqrt{\sigma} \zeta)+k \sqrt{\sigma} \sec (\sqrt{\sigma} \zeta) & , \sigma>0 \\
-k \sqrt{\sigma} \cot (\sqrt{\sigma} \zeta)+k \sqrt{\sigma} \csc (\sqrt{\sigma} \zeta) & , \sigma>0 \\
\frac{2 k}{\zeta} & , \sigma=0
\end{array}\right.
$$

### 3.2 The exact solutions of the first extended (3+1)-dimensional Jimbo-Miwa equation:

$$
\begin{equation*}
u_{x x x y}+3 u_{y} u_{x x}+3 u_{x} u_{x y}+2 u_{y t}-3\left(u_{x z}+u_{y z}+u_{z z}\right)=0 \tag{13}
\end{equation*}
$$

Substituting the wave transformation $u=u(\zeta), \zeta=k x+r y+s z-c t$ into (13) and integrating we get

$$
\begin{equation*}
k^{3} r u^{\prime \prime \prime}+3 k^{2} r\left(u^{\prime}\right)^{2}-\left(3 k s+3 s r+3 s^{2}+2 r c\right) u^{\prime}=0 . \tag{14}
\end{equation*}
$$

Balancing $u^{\prime \prime \prime}$ with $\left(u^{\prime}\right)^{2}$ we get $m=1$. From (4) we have

$$
\begin{equation*}
u(\zeta)=a_{0}+a_{1} \phi(\zeta)+b_{1} \sqrt{\sigma+\phi^{2}(\zeta)} \tag{15}
\end{equation*}
$$

Substituting (5) and (15) into equation (14) and collecting all the coefficients of $\phi^{n}\left(\sqrt{\sigma+\phi^{2}(\zeta)}\right)^{m}$ and setting them equal to zero we obtain

Case 1: $a_{0}=0, \quad a_{1}=-k, \quad b_{1}=k$, and $c=-\left(\frac{k^{3} r \sigma+3 k s+3 s^{2}+3 r s}{2 r}\right)$. The exact solution in this case is

$$
u(\zeta)=\left(\begin{array}{ll}
k \sqrt{-\sigma} \tanh (\sqrt{-\sigma} \zeta)+k \sqrt{\sigma} \operatorname{sech}(\sqrt{-\sigma} \zeta) & , \sigma<0  \tag{16}\\
k \sqrt{-\sigma} \operatorname{coth}(\sqrt{-\sigma} \zeta)-k \sqrt{\sigma} \operatorname{csch}(\sqrt{-\sigma} \zeta) & , \sigma<0 \\
-k \sqrt{\sigma} \tan (\sqrt{\sigma} \zeta)+k \sqrt{\sigma} \sec (\sqrt{\sigma} \zeta) & , \sigma>0 \\
-k \sqrt{\sigma} \cot (\sqrt{\sigma} \zeta)+k \sqrt{\sigma} \csc (\sqrt{\sigma} \zeta) & , \sigma>0 \\
\frac{2 k}{\zeta} & , \sigma=0
\end{array}\right.
$$

Remark: For $\sigma=-1$ the dispersion relation is completely agree with [17]
Case 2: $a_{0}=0, a_{1}=-k, \quad b_{1}=-k, \quad c=-\left(\frac{k^{3} r \sigma+3 k s+3 s^{2}+3 r s}{2 r}\right)$ and the exact solution is

$$
u(\zeta)=\left(\begin{array}{ll}
k \sqrt{-\sigma} \tanh (\sqrt{-\sigma} \zeta)-k \sqrt{\sigma} \operatorname{sech}(\sqrt{-\sigma} \zeta) & , \sigma<0  \tag{17}\\
k \sqrt{-\sigma} \operatorname{coth}(\sqrt{-\sigma} \zeta)+k \sqrt{\sigma} \operatorname{csch}(\sqrt{-\sigma} \zeta) & , \sigma<0 \\
-k \sqrt{\sigma} \tan (\sqrt{\sigma} \zeta)-k \sqrt{\sigma} \sec (\sqrt{\sigma} \zeta) & , \sigma>0 \\
-k \sqrt{\sigma} \cot (\sqrt{\sigma} \zeta)-k \sqrt{\sigma} \csc (\sqrt{\sigma} \zeta) & , \sigma>0 \\
0 & , \sigma=0
\end{array}\right.
$$

Case 3: $a_{0}=0, a_{1}=-2 k, \quad b_{1}=0, \quad c=-\left(\frac{4 k^{3} r \sigma+3 k s+3 s^{2}+3 r s}{2 r}\right)$ and the exact solution is

$$
u(\zeta)=\left(\begin{array}{ll}
2 k \sqrt{-\sigma} \tanh (\sqrt{-\sigma} \zeta) & , \sigma<0  \tag{18}\\
2 k \sqrt{-\sigma} \operatorname{coth}(\sqrt{-\sigma} \zeta) & , \sigma<0 \\
-2 k \sqrt{\sigma} \tan (\sqrt{\sigma} \zeta) & , \sigma>0 \\
-2 k \sqrt{\sigma} \cot (\sqrt{\sigma} \zeta) & , \sigma>0 \\
\frac{2 k}{\zeta} & , \sigma=0
\end{array}\right.
$$

Remark: For $\sigma=-\frac{1}{4}$ the dispersion relation is completely agree with [17]

### 3.3 The exact solutions of the second extended (3+1)-dimensional Jimbo-Miwa equation:

$$
\begin{equation*}
u_{x x x y}+3 u_{y} u_{x x}+3 u_{x} u_{x y}+2\left(u_{x t}+u_{y t}+u_{z t}\right)-3 u_{x z}=0 \tag{19}
\end{equation*}
$$

Substituting the wave transformation $u=u(\zeta), \zeta=k x+r y+s z-c t$ into (19) and integrating we get

$$
\begin{equation*}
k^{3} r u^{\prime \prime \prime}+3 k^{2} r\left(u^{\prime}\right)^{2}-(2 k c+2 s c+2 r c+3 k s) u^{\prime}=0 . \tag{20}
\end{equation*}
$$

Balancing $u^{\prime \prime \prime}$ with $\left(u^{\prime}\right)^{2}$ we get $m=1$. From (4) we have

$$
\begin{equation*}
u(\zeta)=a_{0}+a_{1} \phi(\zeta)+b_{1} \sqrt{\left(\sigma+\phi^{2}(\zeta)\right)} \tag{21}
\end{equation*}
$$

Substituting (5) and (21) into (20) and collecting all the coefficients of $\phi^{n}\left(\sqrt{\sigma+\phi^{2}(\zeta)}\right)^{m}$ and setting them equal to zero we obtain
Case 1: $a_{0}=0, a_{1}=-k, \quad b_{1}=k$, and $c=-\left(\frac{k^{3} r \sigma+3 k s}{2(k+r+s)}\right)$. The exact solution in this case is

$$
u(\zeta)=\left(\begin{array}{ll}
k \sqrt{-\sigma} \tanh (\sqrt{-\sigma} \zeta)+k \sqrt{\sigma} \operatorname{sech}(\sqrt{-\sigma} \zeta) & , \sigma<0,  \tag{22}\\
k \sqrt{-\sigma} \operatorname{coth}(\sqrt{-\sigma} \zeta)-k \sqrt{\sigma} \operatorname{csch}(\sqrt{-\sigma} \zeta) & , \sigma<0, \\
-k \sqrt{\sigma} \tan (\sqrt{\sigma} \zeta)+k \sqrt{\sigma} \sec (\sqrt{\sigma} \zeta) & , \sigma>0, \\
-k \sqrt{\sigma} \cot (\sqrt{\sigma} \zeta)+k \sqrt{\sigma} \csc (\sqrt{\sigma} \zeta) & , \sigma>0, \\
\frac{2 k}{\zeta} & , \sigma=0,
\end{array}\right.
$$

Remark: For $\sigma=-1$ the dispersion relation is completely agree with [17]
Case 2: $a_{0}=0, \quad a_{1}=-k, \quad b_{1}=-k, \quad c=-\left(\frac{k^{3} r \sigma+3 k s}{2(k+r+s)}\right)$ and the exact solution is

$$
u(\zeta)=\left(\begin{array}{ll}
k \sqrt{-\sigma} \tanh (\sqrt{-\sigma} \zeta)-k \sqrt{\sigma} \operatorname{sech}(\sqrt{-\sigma} \zeta) & , \sigma<0  \tag{23}\\
k \sqrt{-\sigma} \operatorname{coth}(\sqrt{-\sigma} \zeta)+k \sqrt{\sigma} \operatorname{csch}(\sqrt{-\sigma} \zeta) & , \sigma<0 \\
-k \sqrt{\sigma} \tan (\sqrt{\sigma} \zeta)-k \sqrt{\sigma} \sec (\sqrt{\sigma} \zeta) & , \sigma>0 \\
-k \sqrt{\sigma} \cot (\sqrt{\sigma} \zeta)-k \sqrt{\sigma} \csc (\sqrt{\sigma} \zeta) & , \sigma>0 \\
0 & , \sigma=0
\end{array}\right.
$$

Case 3: $a_{0}=0, \quad a_{1}=-2 k, \quad b_{1}=0, \quad c=-\left(\frac{4 k^{3} r \sigma+3 k s}{2(k+r+s)}\right)$ and the exact solution is

$$
u(\zeta)=\left(\begin{array}{ll}
2 k \sqrt{-\sigma} \tanh (\sqrt{-\sigma} \zeta) & , \sigma<0  \tag{24}\\
2 k \sqrt{-\sigma} \operatorname{coth}(\sqrt{-\sigma} \zeta) & , \sigma<0 \\
-2 k \sqrt{\sigma} \tan (\sqrt{\sigma} \zeta) & , \sigma>0 \\
-2 k \sqrt{\sigma} \cot (\sqrt{\sigma} \zeta) & , \sigma>0 \\
\frac{2 k}{\zeta} & , \sigma=0
\end{array}\right.
$$

Remark: For $\sigma=-\frac{1}{4}$ the dispersion relation is completely agree with [17]

## 4 Conclusion

In this work, the Riccati mapping equation method has been applied to solve the $(3+1)$ dimensional Jimbo-Miwa equation as well as its two extended forms. We obtained various types of solutions as trigonometric and rational functions. It is verified that these solutions are satisfying their corresponding equations. The obtained results have been compared with Wazwaz [17] and the observation was that our results are new with different solution structures.

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